### UNENE Graduate Course Reactor Thermal-Hydraulics Design and Analysis McMaster University Whitby March 11-12, March 25-26, April 8-9, April 22-23, 2006

# Equation of State

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## Outline

- Thermal-hydraulics Properties
- The Iterative Method
- The Rate Method
- H2O Property Fits

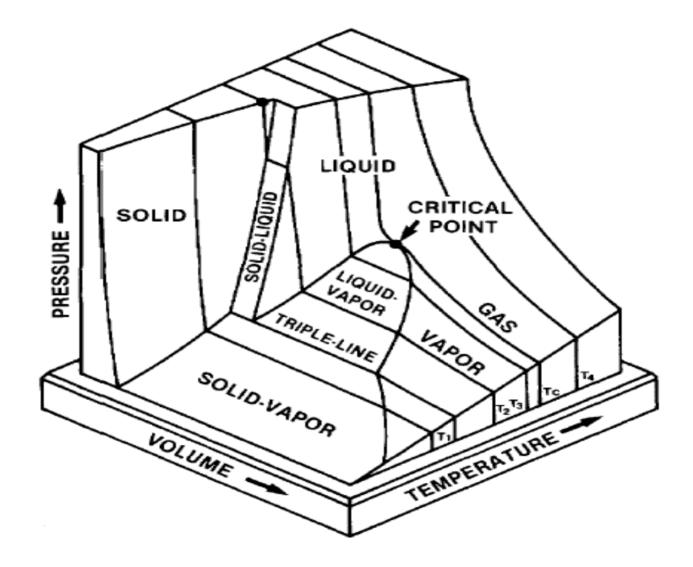
### **Thermo-Hydraulic Properties**

• Equation of state is a relationship between any four thermodynamic properties of a substance

 $\pi$  (P, V, T, M) = 0  $\pi$  (P, v, T) = 0

- Why do we need the equation of state?
  - 4 dependent variables (density, temperature, pressure and velocity); 3 conservation equations (mass, momentum and energy)
- For given values of pressure and temperature, the calculation of other properties is relatively straightforward
- Determination of pressure from known values of other thermodynamic properties requires iteration (characteristic of the equation of state)
- Conservation equations are all cast as rate equations, whereas the equation of state is typically written as an algebraic equation

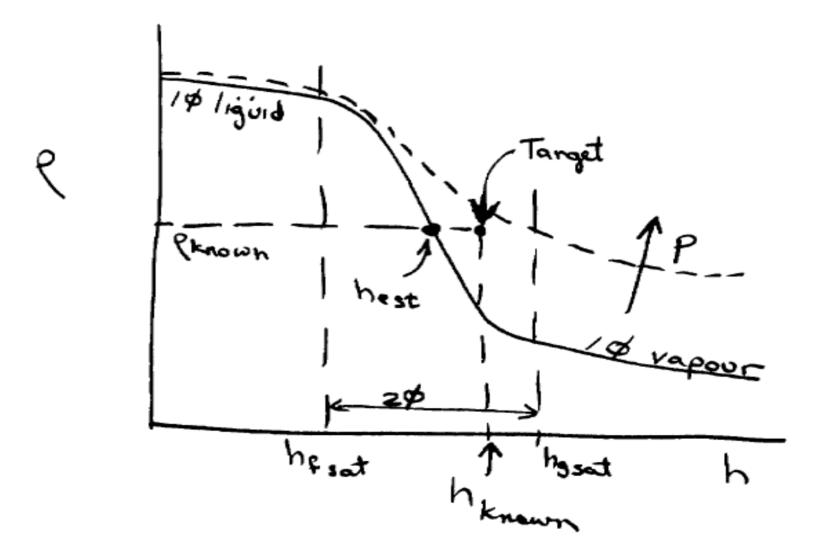
### P-v-T surface of Water



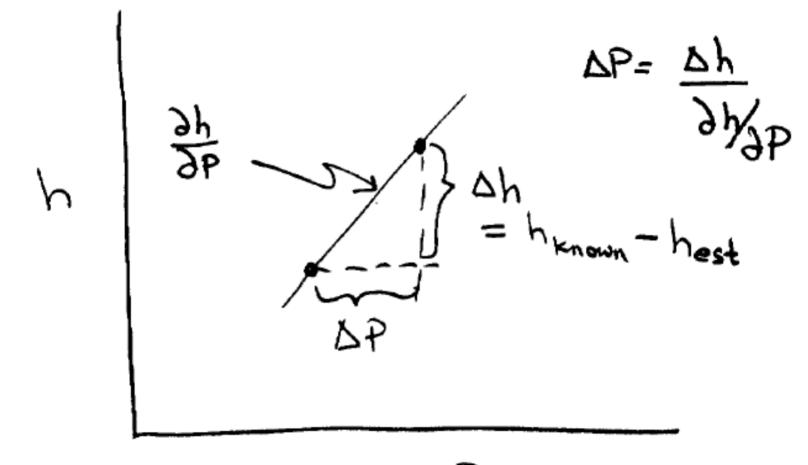
### **The Iterative Method**

- Given the density and enthalpy of a volume of water, the task at hand is to find the associated values of pressure and temperature
- The iteration process
  - Guess the pressure (pressure at previous time step)
  - Calculate  $h_{fsat}$  and  $h_{gsat}$  (saturation enthalpies for gas and liquid)
  - If  $h > h_{gsat}$  the fluid is single-phase steam
  - If  $h \le h_{fsat}$  the fluid is single-phase water
  - Otherwise, the fluid is two-phase mixture

#### Numerical Search for P Given $\rho$ and h



#### Error Correction Scheme



P

#### The Iterative Method – Two-phase Equilibrium

• Given the density and enthalpy of a volume of water, the task at hand is to find the associated values of pressure and temperature

$$v = \frac{1}{\rho} = v_f(P) + x_{est} v_{fg}(P)$$

$$h_{est} = h_f(P) + x_{est} h_{fg}(P)$$

• The pressure iteration continues until the pressure converges with some tolerance

$$\Delta P = \frac{\Delta h}{\left(\partial h/\partial P\right)_{\rho}} \qquad P = P + \Delta P$$

#### **The Iterative Method** – Single-phase

• For single-phase fluid

$$\rho = \rho(P, T)$$
 and  $h = h(P, T)$ 

• For guessed pressure and temperature, density and enthalpy can be found directly from property tables; Taylor's series expansion is used:

$$\rho = \rho_{est} + \frac{\partial \rho}{\partial T} \bigg|_{P} \Delta T + \frac{\partial \rho}{\partial P} \bigg|_{T} \Delta P \qquad h = h_{est} + \frac{\partial h}{\partial T} \bigg|_{P} \Delta T + \frac{\partial h}{\partial P} \bigg|_{T} \Delta P$$

• Using a guess for the increment of density and enthalpy, the increments for pressure and temperature can be found

$$\Delta \rho = \rho - \rho_{est} \qquad \Delta h = h - h_{est}$$

#### **The Iterative Method** – Single-phase

$$\Delta P = \frac{\frac{\partial h}{\partial T}}{\frac{\partial \rho}{\partial P}_{T} \frac{dh}{dT}_{P}} \Delta \rho - \frac{\partial \rho}{\partial T}_{P} \Delta h \qquad \Delta T = \frac{\frac{\partial h}{\partial P}}{\frac{\partial \rho}{\partial T}_{T} \Delta \rho - \frac{\partial \rho}{\partial T}_{T}} \Delta h \qquad \Delta T = \frac{\frac{\partial h}{\partial P}}{\frac{\partial \rho}{\partial T}_{P} \frac{dh}{dP}_{T}} - \frac{\frac{\partial \rho}{\partial T}}{\frac{\partial \rho}{\partial T}_{P} \frac{dh}{dT}_{P}}$$

$$\Delta P = G_{1P} \Delta \rho + G_{2P} \Delta h$$

 $\Delta T = G_{1T} \ \Delta \rho + G_{2T} \ \Delta h$ 

 $P = P + \Delta P$  and  $T = T + \Delta T$ 

• The iteration is repeated until until pressure and temperature converge with some tolerance

### The Rate Method

- The rate method eliminates the need for iteration with no loss of accuracy
- The rate method is consistent with the conservation equation form
- The current form of the rate equation consistent with the Porsching Method

• We wish to relate the rate of change of pressure to rates of change in density and enthalpy:

$$dP = G_1 d\rho + G_2 dh \qquad \frac{dP}{dt} = G_1 \frac{d\rho}{dt} + G_2 \frac{dh}{dt}$$

• For two-phase homogeneous mixture, and differentiating the equations below:

$$\mathbf{v} = \mathbf{v}_{f} + \mathbf{x}\mathbf{v}_{fg} \qquad \mathbf{v}_{fg} \equiv \mathbf{v}_{g} - \mathbf{v}_{f} \text{ and } \mathbf{h}_{fg} \equiv \mathbf{h}_{g} - \mathbf{h}_{f},$$

$$\mathbf{h} = \mathbf{h}_{f} + \mathbf{x}\mathbf{h}_{fg}$$

$$\frac{d\mathbf{h}}{d\mathbf{t}} = \left(\frac{\partial \mathbf{h}}{\partial \mathbf{P}}\right)_{\rho} \frac{d\mathbf{P}}{d\mathbf{t}} = \left[\frac{\partial \mathbf{h}_{f}}{\partial \mathbf{P}} + \mathbf{h}_{fg} \frac{\partial \mathbf{x}}{\partial \mathbf{P}} + \mathbf{x} \frac{\partial \mathbf{h}_{fg}}{\partial \mathbf{P}}\right] \frac{d\mathbf{P}}{d\mathbf{t}}.$$

$$\frac{d\mathbf{x}}{d\mathbf{P}} = \frac{\partial \left(\frac{\mathbf{v} - \mathbf{v}_{f}}{\mathbf{v}_{fg}}\right)}{\partial \mathbf{P}} = -\frac{1}{\mathbf{v}_{fg}} \left[\frac{\partial \mathbf{v}_{f}}{\partial \mathbf{P}} + \mathbf{x} \frac{\partial \mathbf{v}_{fg}}{\partial \mathbf{P}}\right].$$
 for density constant

$$\begin{aligned} \frac{dh}{dt} &= \left\{ \frac{\partial h_{f}}{\partial P} + x \frac{\partial h_{fg}}{\partial P} - \frac{h_{fg}}{v_{fg}} \left[ \frac{\partial v_{f}}{\partial P} + x \frac{\partial v_{fg}}{\partial P} \right] \right\} \frac{dP}{dt} \\ \frac{dP}{dt} &= \frac{v_{fg}}{\left\{ v_{fg} \left[ \frac{\partial h_{f}}{\partial P} + x \frac{\partial h_{fg}}{\partial P} \right] - h_{fg} \left[ \frac{\partial v_{f}}{\partial P} + x \frac{\partial v_{fg}}{\partial P} \right] \right\}}{\left\{ v_{fg} \left[ \frac{\partial h_{f}}{\partial P} + x \frac{\partial h_{fg}}{\partial P} \right] - h_{fg} \left[ \frac{\partial v_{f}}{\partial P} + x \frac{\partial v_{fg}}{\partial P} \right] \right\}} \frac{dh}{dt} \\ &= \frac{v_{fg}}{\left\{ \text{DENOMINATOR} \right\}} \frac{dh}{dt} = G_{2} \frac{dh}{dt} . \end{aligned}$$
$$\begin{aligned} \frac{dP}{dt} &= \frac{-h_{fg}}{\left\{ \text{DENOMINATOR} \right\}} \frac{dv}{dt} = \frac{h_{fg} v^{2}}{\left\{ \text{DENOMINATOR} \right\}} \frac{d\rho}{dt} = G_{1} \frac{d\rho}{dt} . \end{aligned}$$

$$p = M/V \text{ and } h = H/M,$$

$$\frac{d\rho}{dt} = \frac{1}{V} \frac{dM}{dt} - \frac{M}{V^2} \frac{dV}{dt}$$

$$\frac{dh}{dt} = \frac{1}{M} \frac{dH}{dt} - \frac{H}{M^2} \frac{dM}{dt}.$$

$$\frac{dP}{dt} = \left(\frac{G_1}{V} - \frac{G_2H}{M^2}\right) \frac{dM}{dt} + \frac{G_2}{M} \frac{dH}{dt} - \frac{G_1M}{V^2} \frac{dV}{dt}.$$

$$\frac{dP}{dt} = \frac{F_1 \frac{dM}{dt} + F_2 \frac{dH}{dt} + F_3 \frac{dV}{dt}}{M_g F_4 + M_f F_5}$$

$$\begin{split} F_{1} &= h_{g} \nu_{f} - h_{f} \nu_{g} \\ F_{2} &= \nu_{g} - \nu_{f} \\ F_{3} &= h_{f} - h_{g} \\ F_{4} &= \frac{\partial h_{g}}{\partial P} (\nu_{g} - \nu_{f}) - \frac{\partial \nu_{g}}{\partial P} (h_{g} - h_{f}) \\ F_{5} &= \frac{\partial h_{f}}{\partial P} (\nu_{g} - \nu_{f}) - \frac{\partial \nu_{f}}{\partial P} (h_{g} - h_{f}) \\ M_{g} &\equiv x M \\ M_{f} &\equiv (1 \ x) M . \end{split}$$

• We wish to relate the rate of change of pressure to rates of change in density and enthalpy:

$$\frac{\mathrm{d}P}{\mathrm{d}t} = G_{1P} \frac{\mathrm{d}\rho}{\mathrm{d}t} + G_{2P} \frac{\mathrm{d}h}{\mathrm{d}t} \qquad \frac{\mathrm{d}T}{\mathrm{d}t} = G_{1T} \frac{\mathrm{d}\rho}{\mathrm{d}t} + G_{2T} \frac{\mathrm{d}h}{\mathrm{d}t} \,.$$

• For single-phase fluid we do not need to account for the sotring out between the phases, thus the derivation is more direct and less complex

$$\mathbf{P} = \pi \left( \rho, \mathbf{h} \right) \qquad \qquad \rho = \rho \left( \mathbf{P}, \mathbf{T} \right)$$

$$\frac{dP}{dt} = \frac{\partial P}{\partial \rho} \int_{h} \frac{d\rho}{dt} + \frac{\partial P}{\partial h} \int_{\rho} \frac{dh}{dt} \qquad \qquad h = h(P, T)$$

$$\frac{d\rho}{dt} = \frac{\partial\rho}{\partial P} \bigg|_{T} \frac{dP}{dt} + \frac{\partial\rho}{\partial T} \bigg|_{P} \frac{dT}{dt}$$

$$\frac{\mathrm{d}\mathbf{h}}{\mathrm{d}\mathbf{t}} = \frac{\partial \mathbf{h}}{\partial \mathbf{P}} \bigg|_{\mathbf{T}} \frac{\mathrm{d}\mathbf{P}}{\mathrm{d}\mathbf{t}} + \frac{\partial \mathbf{h}}{\partial \mathbf{T}} \bigg|_{\mathbf{P}} \frac{\mathrm{d}\mathbf{T}}{\mathrm{d}\mathbf{t}} \, .$$

• Solving for dP/dt and dT/dt we get:

$$\frac{dP}{dt} = \frac{\frac{\partial h}{\partial T}}{\frac{\partial \rho}{\partial P}_{T} \frac{dh}{dT}_{P} \frac{d\rho}{dt}} - \frac{\partial \rho}{\partial T}_{P} \frac{dh}{dt}_{P} \frac{dh}{dt} = \frac{\frac{\partial h}{\partial P}}{\frac{\partial \rho}{\partial T}_{T} \frac{d\rho}{dt}} - \frac{\partial \rho}{\partial T}_{T} \frac{dh}{dt}_{T} \frac{dh}{dt}$$

• The extensive form of the equations is also easily obtained

$$\frac{dP}{dt} = \frac{F_{1P} \frac{dM}{dt} + F_{2P} \frac{dH}{dt} + F_{3P} \frac{dV}{dt}}{M_v F_{4P} + M_\ell F_{5P}}$$

$$\frac{dT}{dt} = \frac{F_{1T}}{\frac{dM}{dt}} + F_{2T}\frac{dH}{dt} + F_{3T}\frac{dV}{dt}}{M_vF_{4T}} + M_\ell F_{5T}$$

$$\begin{split} F_{1P} &= \rho \left( \frac{\partial h}{\partial T} \right)_{P} + h \left( \frac{\partial \rho}{\partial T} \right)_{P} \\ F_{2P} &= - \frac{\partial \rho}{\partial t} \right)_{P} \\ F_{3P} &= -\rho^{2} \left( \frac{\partial \rho}{\partial T} \right)_{P} \\ F_{4P} &= 0 \text{ for subcooled, } = \left( \frac{d\rho}{\partial P} \right)_{T} \left( \frac{\partial h}{\partial T} \right)_{P} - \left( \frac{\partial \rho}{\partial T} \right)_{P} \left( \frac{dh}{dP} \right)_{T} \text{ for superheated} \\ F_{5P} &= \left( \frac{d\rho}{\partial P} \right)_{T} \left( \frac{\partial h}{\partial T} \right)_{P} - \left( \frac{\partial \rho}{\partial T} \right)_{P} \left( \frac{dh}{dP} \right)_{T} \text{ for subcooled} = 0 \text{ for superheated} \\ F_{1T} &= \rho \left( \frac{\partial h}{\partial T} \right)_{P} + h \left( \frac{\partial \rho}{\partial T} \right)_{P} \\ F_{2T} &= - \left( \frac{\partial \rho}{\partial t} \right)_{T} \\ F_{3T} &= -\rho \left( \frac{\partial h}{\partial T} \right)_{T} \\ F_{4T} &= -F_{4P} \\ F_{5T} &= -F_{5P} \\ M_{v} &= \text{mass of vapour phase} = 0 \text{ for subcooled, } = M \text{ for superheated} \\ M_{f} &= \text{mass of liquid phase} = M \text{ for subcooled} = 0 \text{ for superheated} \end{split}$$

• The rate form for the two-phase non-equilibrium case is a simple extension of the single-phase non-equilibrium case

$$\frac{dP_k}{dt} = G_{1P}^k \frac{d\rho_k}{dt} + G_{2P}^k \frac{dh_k}{dt}$$
$$\frac{dT_k}{dt} = G_{1T}^k \frac{d\rho_k}{dt} + G_{2T}^k \frac{dh_k}{dt}$$

Where k represents either water or steam

• The derivates of density and enthalpy are available in the rate form of the equation of state

#### The Rate Method - Summary of the F Functions

Case	F <sub>1</sub>	F <sub>2</sub>	F <sub>3</sub>	$F_4$	F <sub>5</sub>
2φ equilibrium (all derivatives along saturation line)	$h_g v_f - h_f v_g$	ν <sub>fg</sub>	h <sub>fg</sub>	$\frac{\partial h_g}{\partial P} v_{fg} - \frac{\partial v_g}{\partial P} h_{fg}$	$\frac{\partial  h_f}{\partial  P}   \nu_{fg}  -  \frac{\partial  \nu_f}{\partial  P}   h_{fg}$
lφ non- equilibrium pressure	$ \rho \left( \frac{\partial h}{\partial T} \right)_{p} + h \left( \frac{\partial \rho}{\partial T} \right)_{p} $	$-\frac{\partial \rho}{\partial T}\Big _{P}$	$-\rho^2 \left( \frac{\partial h}{\partial T} \right)_p$	$\frac{\partial \rho}{\partial P} \int_{T} \frac{\partial h}{\partial T} \int_{P} - \frac{\partial \rho}{\partial T} \int_{P} \frac{\partial h}{\partial P} \int_{T} \frac{\partial h}{\partial P}$ superheated	$\left(\frac{\partial \mathbf{p}}{\partial \mathbf{r}}\right)_{\mathrm{T}} \left(\frac{\partial \mathbf{h}}{\partial \mathbf{T}}\right)_{\mathrm{P}} = \left(\frac{\partial \mathbf{p}}{\partial \mathbf{T}}\right)_{\mathrm{P}} \left(\frac{\partial \mathbf{h}}{\partial \mathbf{P}}\right)_{\mathrm{T}}$ subcooled 0 superheated
lφ non- equilibrium temperature	$\left. \rho \; \frac{\partial h}{\partial P} \right)_T + h \; \frac{\partial \rho}{\partial P} \right)_T$	$-\frac{\partial \rho}{\partial P}$	$\left\rho^2 \; \frac{\partial h}{\partial P}\right)_T$	$-F_{4P}$	– F <sub>5P</sub>

Table 4.1 Summary of the F functions for the rate form of the equation of state

### The Rate Method - Summary of the G Functions

Table 4.2 Summary of the G functions for the rate form of the equation of state

Case	Gl	G <sub>2</sub>
2φ equilibrium (all derivatives along saturation line)	$\frac{h_{fg} v^2}{\left\{ v_{fg} \left[ \frac{\partial h_f}{\partial P} + x \; \frac{\partial h_{fg}}{\partial P} \right] - h_{fg} \left[ \frac{\partial v_f}{\partial P} + x \; \frac{\partial v_{fg}}{\partial P} \right] \right\}}$	$\frac{\nu_{fg}}{\left\{\nu_{fg}\left[\frac{\partial h_{f}}{\partial P} + x \frac{\partial h_{fg}}{\partial P}\right] - h_{fg}\left[\frac{\partial \nu_{f}}{\partial P} + x \frac{\partial \nu_{fg}}{\partial P}\right]\right\}}$
1φ non-equilibrium pressure	$\frac{\frac{\partial h}{\partial T}}{\left[\frac{\partial \rho}{\partial P}\right]_{T} \frac{\partial h}{\partial T}}_{P} = \frac{\partial \rho}{\partial T}_{P} \frac{\partial h}{\partial P}_{T}$	$ -\frac{\partial \rho}{\partial T} \bigg _{P} $ $ \left[ \frac{\partial \rho}{\partial P} \bigg _{T} \frac{\partial h}{\partial T} \bigg _{P} - \frac{\partial \rho}{\partial T} \bigg _{P} \frac{\partial h}{\partial P} \bigg _{T} \right] $
1φ non-equilibrium temperature	$\frac{\frac{\partial h}{\partial P}}{\frac{\partial r}{\partial T}}_{P} \frac{\frac{\partial h}{\partial P}}{\frac{\partial P}{T}}_{T} - \frac{\frac{\partial \rho}{\partial P}}{\frac{\partial P}{T}}_{T} \frac{\frac{\partial h}{\partial T}}{\frac{\partial T}{P}}_{P}$	$-\frac{\partial \rho}{\partial P}\Big _{T}$ $\left[\frac{\partial \rho}{\partial T}\Big _{P}\frac{\partial h}{\partial P}\Big _{T}-\frac{\partial \rho}{\partial P}\Big _{T}\frac{\partial h}{\partial T}\Big _{P}\right]$

# H2O Property Fits

- Steam-water properties are are calculated using standard routines
  - CATHENA has built in and stand-alone property tables
  - TH codes use different forms of property tables need to make sure understand differences in property tables that can create some differences in results
- The overall approach is based on careful attention to the fitting around the saturation line which separates single-phase from two-phase applications
  - Most codes operate near saturation line

# H2O Property Fits

• Density, enthalpy and other properties vary more strongly with temperature than with pressure

$$\rho(\mathbf{P},\mathbf{T}) = \rho_{sat}(\mathbf{T}) + \frac{\partial \rho}{\partial \mathbf{P}} \int_{\mathbf{T}} (\mathbf{P} - \mathbf{P}_{sat}(\mathbf{T}))$$

- Both the properties and their derivatives must be free of discontinuities to establish numerical stability
- Discontinuities in properties can arise especially due to flow regime changes
  - Therefore, carefully chosen smoothing techniques must be used to avoid sudden changes in properties due to flow regime changes

### Density vs Pressure Diagram

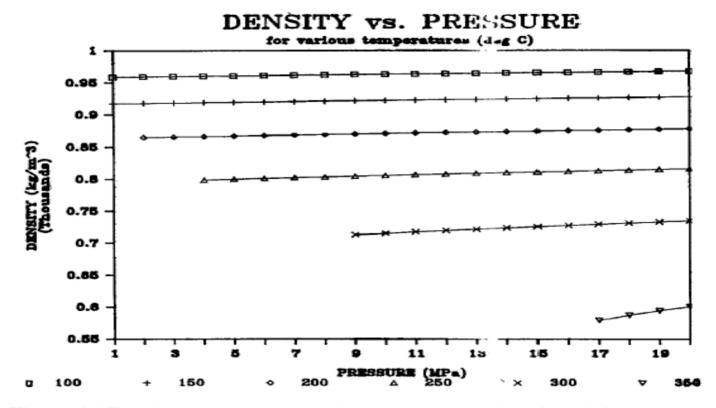


Figure 4.4 Density vs. pressure at various temperatures in subcooled water.

