

#### 4. Computational Scheme for CANDU Neutronics

- The computational scheme for CANDU neutronics consists of three stages:
- Cell calculation: to determine lattice properties for basic lattice cells
- "Supercell" calculation: to determine the "incremental" cross sections to be added to the basic-cell properties to account for the effect of reactivity devices
- Finite-core calculation: to solve the neutron-diffusion problem in the reactor core, to calculate the 3dimensional flux and power distribution.



#### 4. Computational Scheme for CANDU Neutronics

- Computer programs have been developed to perform the calculations corresponding to each stage in the above process.
- These are now briefly discussed in the following sections.



- The cell calculation treats the "bare" CANDU basic lattice cell.
- "B are" here means the basic lattice cell without reactivity devices superimposed (refer to Fig. 1.2).



- The cell (or lattice) code which has traditionally been used for CANDU design and analysis is POWDERPUFS-V.
- This is an empirical-recipe code, based on the results of measurements made on heavy-water-moderated lattices in research reactors ZEEP and ZED-2 at Chalk River Laboratories.
- Although based on empiricism rather than a strong theoretical foundation, POWDERPUFS-V has been applied very successfully to CANDU design and analysis.



- POWDERPUFS-V uses the four-factor formula to calculate the infinite-lattice multiplication constant k<sub>inf</sub> = εpηf,
- and to calculate homogenized-cell nuclear cross sections.
- It also utilizes the Westcott formulation for nuclide cross sections, a parametrization in terms of the neutron temperature and a spectral parameter r.



- The Westcott parametrization is applicable to highly thermalized neutron spectra,
- such as those in the CANDU lattice cell,
- where over 95% of neutrons in the fuel have a Maxwellian energy distribution.



- The nuclear cross sections are evaluated using the Westcott formula and other simple recipes,
- using parameter values obtained empirically from experiment.
- POWDERPUFS-V is applicable to CANDU reactors fuelled with natural uranium,
- where the amount of plutonium in the fuel is limited by the natural-uranium burnup.



- The Westcott convention for calculating the effective cross sections of fuel nuclides is based on assuming
- that the neutron spectrum can be written as the sum of a Maxwellian function and an epithermal function tending to 1/E:

 $n(v) = N(1-f)\rho_m(v) + Nf\rho_e(v)$  (4.1) where  $\rho_m(v)$  and  $\rho_e(v)$  are the Maxwellian and epithermal normalized density distribution functions, respectively, and



N = total neutron density f = fraction of the total neutron density in the epithermal spectrum

$$\rho_{\rm m}({\bf v}) = \frac{4}{\sqrt{\pi}} \frac{{\bf v}^2}{{\bf v}_{\rm T}^3} e^{-({\bf v}/{\bf v}_{\rm T})^2}$$
(4.2)  
$$\rho_{\rm e}({\bf v}) = {\bf v}_{\rm T} \sqrt{\mu} \frac{\Delta({\bf v})}{{\bf v}^2}$$
(4.3)

and  $v_T$  = velocity of a neutron of energy kT



- Δ(v) is an empirical function describing the way the epithermal spectrum (with its 1/E "tail") joins the Maxwellian spectrum.
- It satisfies  $\Delta(\mathbf{v}) \rightarrow 0$  for  $\mathbf{E} < \mu \mathbf{k} \mathbf{T}$
- and  $\Delta(\mathbf{v}) \rightarrow 1$  for  $\mathbf{E} > \mu \mathbf{kT}$ ,
- where μkT represents the lower limit of the 1/E spectrum and, by choice of convention, μ = 3.681.



• The Westcott flux  $\hat{\phi}$  is defined as  $\hat{\phi} = Nv_0$ where  $v_0 = 2200$  m/s

(i.e., as if the entire neutron distribution had a speed of 2200 m/s)



- + and the effective (Westcott) cross section  $\hat{\sigma}$
- of a given nuclide is defined so that its product with the Westcott flux gives the total reaction rate: Total reaction rate in nuclide

$$=\hat{\sigma}\hat{\phi}=\hat{\sigma}N \mathbf{v}_{0}$$
(4.5)

• (By total reaction rate is meant the reaction rate in the entire spectrum, which includes the Maxwellian and the 1/E parts.)



- It can be shown that  $\hat{\sigma}$  can be written in terms of 0, the 2200-m/s cross section, as follows:  $\hat{\sigma} = \sigma_0(g + rs)$  (4.6)
- where g is the ratio of the reaction rate of the nuclide in a pure Maxwellian spectrum to the reaction rate of a 1/v absorber of the same 2200 m/s cross section
- (i.e., g is a measure of the 'non-1/v' character of the absorber in a Maxwellian spectrum), and
- r is a measure of the epithermal part (i.e., the 'hardness' of the spectrum).



- r has a small value in the CANDU lattice:
- typically, r ~ 0.05 for a CANDU lattice fuelled with natural UO<sub>2</sub>.
- This is what makes the Westcott formulation a good approximation in CANDU reactors.



- The g and s values for the various nuclides are obtained from experiment
- They are functions of the neutron temperature T<sub>n</sub>, so that Eq. (4.6) is evaluated in fact as

 $\hat{\sigma}(\mathbf{r}, \mathbf{Tn}) = \sigma_0(\mathbf{g}(\mathbf{Tn}) + \mathbf{rs}(\mathbf{Tn}))$  (4.7)



- In POWDERPUFS-V, the factors g(Tn) and s(Tn) are expressed as power series in the neutron temperature.
- With this database of g and s values for various nuclides, POWDERPUFS-V can calculate reaction rates in the fuel very quickly, using Eq. (4.7).
- For other materials, constant inputs or simple recipes are used.



- The methodology requires the evaluation of the spectral parameter r and the neutron temperature T<sub>n</sub>
- in an iterative fashion from the lattice parameters via empirical relationships.



- In addition to the empiricisms, there are some approximations:
- Fast fission is taken into account in <sup>238</sup>U only, and is "lumped" into the thermal-fission cross section;
- also, up-scattering is ignored.



- One great advantage of POWDERPUFS-V is that, due to its semi-empirical nature and the simplifying assumptions used, it is very fast-running:
- A complete calculation (including depletion to exit burnup values) for a given lattice type is performed in less than 1 second of CPU time.



- POWDERPUFS-V provides "homogenized-cell" twoenergy-group lattice properties for input into finitecore models and calculations.
- In standard "fuel-burn" mode, the lattice properties are provided as functions of fuel irradiation (or burnup)
- for specified values of lattice conditions, such as geometry, fuel, coolant, and moderator temperatures, power level, coolant density, moderator-poison concentration, etc.
- These conditions are entered as input to the code.



- For instance, the geometrical input and fuel mass may correspond to those for the 37-element-natural fuel lattice,
- the (average) fuel temperature may be entered as 687° C,
- the (average) coolant temperature 290° C,
- the moderator temperature 70° C,
- the moderator and coolant purities may be 99.9 and 99.75 weight % D<sub>2</sub>O, and
- the moderator poison may be set to 0 ppm B



- In addition to providing nuclear properties at various values of instantaneous irradiation,
- **POWDERPUFS-V** can do a "reaction-rateaveraged" calculation,
- where it averages the properties over irradiation ω from 0 to a specified exit value:

$$\Sigma_{r.r.av.} = \left(\frac{1}{\omega_{exit}}\right) \int_{0}^{\omega_{exit}} \Sigma(\omega) d\omega$$

• This calculation is extremely useful as a fairly good point-reactor model.



- A "perturbation" mode is also provided in POWDERPUFS-V,
- where lattice properties are evaluated assuming "instantaneous" changes in lattice conditions
- occurring at various values of fuel irradiation.



- POWDERPUFS-V has been used as the lattice code for CANDU reactors for about 30 years,
- where it has performed very well.
- One advantage of POWDERPUFS-V is that it is incorporated as a module within the finite-core code RFSP, described in Section 4.3 below.



- While POWDERPUFS-V has traditionally been the lattice code for CANDU design and analysis,
- it will eventually be replaced by a code with a stronger theoretical foundation, a multigroup transport-theory code such as WIMS-AECL.
- Calculating lattice cross sections with WIMS-AECL is, however, more complex and
- much more computationally intensive than with **POWDERPUFS-V.**



- The effects of reactivity devices on the nuclear properties of the lattice in their vicinity are determined by a supercell calculation,
- performed with the computer code MULTICELL.



- A typical supercell is shown in Fig. 4.1.
- It is essentially a small model volume of the core around a portion of the reactivity device,
- including a portion of the neighbouring fuel channel (normally oriented perpendicularly to the device).
- The dimensions of the supercell are typically 1 lattice pitch x 0.5 lattice pitch x 0.5 bundle length.
- This represents a unit volume over which the effect of the reactivity device is modelled, utilizing the assumption of mirror symmetry about the supercell boundaries.



- The calculation provides incremental cross sections, which are to be added to the basic lattice cross sections over "homogenized" supercell volumes along the length of the device.
- MULTICELL applies pre-calculated boundary conditions (current-to-flux ratios) on internal surfaces which represent the reactivity device and the fuel (modified to Cartesian geometry). These boundary conditions are calculated using integral transport theory (Kushneriuk's method).



- Outside the reactivity-device and fuel regions, MULTICELL applies diffusion theory to calculate the 3-dimensional flux distribution in the moderator.
- Except for zone controllers, which have a spectral effect, 1-group instead of 2-group diffusion is used, and the thermal-neutron source distribution is assumed flat in the moderator,
- based on the contribution to the slowing-down density of the multiple line sources represented by the fuel channels.



- Once the flux distribution in the supercell has been calculated, MULTICELL derives the homogenizedsupercell cross sections.
- Then, from two MULTICELL calculations performed for the supercell:
  - a reference calculation with the device absent, and
  - a calculation with the device present
- device incremental cross sections are obtained by subtraction



- Once basic-lattice properties and reactivity-device incremental cross sections are available,
- the finite-core calculation can proceed.
- The finite-core computer code RFSP (Reactor Fuelling Simulation Program) is specifically designed for CANDU reactors.



- It can calculate the steady-state 3-dimensional flux and power distributions in the reactor using two different methods:
- by solving the time-independent finite-difference diffusion equation in two energy groups, and
- by the method of flux mapping (described in Section 3.3), if the readings of the in-core vanadium detectors are available.



 The time-independent neutron-diffusion equation solved in RFSP for eigenvalue problems in two energy groups with lattice properties from POWDERPUFS-V is:

$$-\vec{\nabla} \cdot \mathbf{D}_{1}(\vec{r}) \vec{\nabla} \phi_{1}(\vec{r}) + (\Sigma_{a1}(\vec{r}) + \Sigma_{1 \to 2}(\vec{r})) \phi_{1}(\vec{r}) - \frac{\nu \Sigma_{f}(\vec{r})}{k_{eff}} \phi_{2}(\vec{r}) = \mathbf{0}$$

 $-\vec{\nabla} \cdot \mathbf{D}_{2}(\vec{\mathbf{r}}) \vec{\nabla} \phi_{2}(\vec{\mathbf{r}}) + \Sigma_{a2}(\vec{\mathbf{r}}) \phi_{2}(\vec{\mathbf{r}}) - \Sigma_{1 \to 2}(\vec{\mathbf{r}}) \phi_{1}(\vec{\mathbf{r}}) = \mathbf{0}$ 



- In this equation there are no fast-fission or
- up-scattering terms, consistent with the POWDERPUFS-V methodology
- However, for use with WIMS-AECL lattice properties, RFSP has the capability to solve a truetwo-energy-group diffusion equation::

$$-\vec{\nabla} \cdot \mathbf{D}_{1}(\vec{\mathbf{r}}) \,\vec{\nabla} \phi_{1}(\vec{\mathbf{r}}) + (\Sigma_{a1}(\vec{\mathbf{r}}) + \Sigma_{1 \rightarrow 2}(\vec{\mathbf{r}})) \,\phi_{1}(\vec{\mathbf{r}}) - \left(\Sigma_{2 \rightarrow 1}(\vec{\mathbf{r}}) + \frac{\nu \Sigma_{f2}(\vec{\mathbf{r}})}{k_{eff}}\right) \phi_{2}(\vec{\mathbf{r}}) = \mathbf{0}$$

$$-\vec{\nabla} \cdot \mathbf{D}_{2}(\vec{\mathbf{r}}) \,\vec{\nabla} \phi_{2}(\vec{\mathbf{r}}) + \left( \Sigma_{a2}(\vec{r}) + \Sigma_{2 \to 1}(\vec{r}) \right) \phi_{2}(\vec{\mathbf{r}}) - \left( \Sigma_{1 \to 2}(\vec{r}) + \frac{\nu \Sigma_{f1}(\vec{r})}{k_{eff}} \right) \phi_{1}(\vec{\mathbf{r}}) = \mathbf{0}$$



• A typical reactor model used with RFSP is shown in Fig. 4.2 a and b (face and top views respectively).



- Major applications of RFSP are in:
- core-design calculations and analyses, including fuelmanagement design calculations, and simulations of reactor power histories
- core-follow calculations at CANDU sites, to track the actual reactor operating history, with burnup steps and channel refuellings.



- Additional capabilities of the program include, among others:
- the calculation of flux distributions for various reactor configurations
- the simulation of <sup>135</sup>Xe/<sup>135</sup>I transients
- the capability for simulating (quasi-statically) bulk control and spatial control
- the calculation of harmonic flux shapes for use in flux mapping,





- the calculation of the reactivity increase expected on refuelling of individual fuel channels
- the capability for solving neutron-kinetics problems by the Improved Quasi-Static (IQS) method.
- RFSP can therefore be used to analyze fast transients, such as those following hypothetical large-loss-of-coolant accidents (LOCA), and can be used to simulate and verify the performance of the shutdown systems.